

Generalized Cesàro operators: geometry of spectra and quasi-nilpotency

Adem Limani, Bartosz Malman

Abstract

For the class of Hardy spaces and standard weighted Bergman spaces of the unit disk we prove that the spectrum of a generalized Cesàro operator T_g is unchanged if the symbol g is perturbed to $g + h$ by an analytic function h inducing a quasi-nilpotent operator T_h , i.e. spectrum of T_h equals $\{0\}$. We also show that any T_g operator which can be approximated in the operator norm by an operator T_h with bounded symbol h is quasi-nilpotent. In the converse direction, we establish an equivalent condition for the function $g \in \mathbf{BMOA}$ to be in the \mathbf{BMOA} -norm closure of H^∞ . This condition turns out to be equivalent to quasi-nilpotency of the operator T_g on the Hardy spaces. This raises the question whether similar statement is true in the context of Bergman spaces and the Bloch space. Furthermore, we provide some general geometric properties of the spectrum of T_g operators.

I INTRODUCTION

Let \mathbb{D} denote the unit disk of the complex plane \mathbb{C} . The generalized Cesàro operator with analytic symbol $g : \mathbb{D} \rightarrow \mathbb{C}$ acts on an analytic function $f : \mathbb{D} \rightarrow \mathbb{C}$ by

$$T_g f(z) = \int_0^z f(\zeta) g'(\zeta) d\zeta, \quad z \in \mathbb{D}.$$

We will for the most part be working in the context of T_g acting on the Hardy spaces and the standard weighted Bergman spaces. It is a classical result in the theory that T_g is bounded on the Hardy spaces H^p , for $0 < p < \infty$, if and only if the symbol g lies in the space \mathbf{BMOA} (see [6]). The Hardy spaces H^p are defined, as usual, to be the spaces of analytic functions in \mathbb{D} which satisfy

$$\|f\|_{H^p}^p := \sup_{0 < r < 1} \int_{\mathbb{T}} |f(r\zeta)|^p dm(\zeta) = \int_{\mathbb{T}} |f(\zeta)|^p dm(\zeta) < \infty,$$

where dm denotes the normalized Lebesgue measure on the unit circle $\mathbb{T} = \{z : |z| = 1\}$ and the last integral is defined by the boundary values of f . The space \mathbf{BMOA} is the dual

of H^1 under the usual Cauchy pairing

$$\lim_{r \rightarrow 1^-} \int_{\mathbb{T}} f(\zeta) \overline{g(r\zeta)} dm(\zeta), \quad f \in H^1, g \in \mathbf{BMOA}.$$

The weighted Bergman space $L_a^{p,\alpha}$, for $0 < p < \infty$ and $\alpha > -1$, consists of analytic functions in \mathbb{D} which satisfy

$$\|f\|_{L^{p,\alpha}}^p := \int_{\mathbb{D}} |f(z)|^p (1 - |z|^2)^\alpha dA(z) < \infty,$$

where dA denotes the normalized area measure. In this case, the criterion for boundedness of T_g is that the symbol g belongs to the Bloch space \mathcal{B} (see, for instance, [2]), which is the space of analytic functions in \mathbb{D} satisfying

$$\sup_{z \in \mathbb{D}} (1 - |z|^2) |g'(z)| < \infty.$$

The purpose of the present paper is to shed light on some spectral properties of T_g operators, with special emphasis on the quasi-nilpotent operators.

We start by briefly describing a common approach that has been followed with the purpose of characterizing the spectrum of T_g operators for several classes of Banach spaces X of analytic functions. We will always implicitly assume that the evaluations $f \mapsto f(\lambda)$ are bounded on any space considered. Since $T_g f(0) = 0$ for any g and f , if X contains functions which do not vanish at $z = 0$, then T_g is not surjective, and thus the point 0 is always contained in the spectrum. It is also easy to see that the operator T_g has no eigenvalues, and that for any $\lambda \in \mathbb{C} \setminus \{0\}$ and any analytic h , the equation $(I - \lambda^{-1} T_g) f = h$ has the unique solution

$$f(z) = R_g(\lambda) h(z) := h(0) e^{\frac{g(z)}{\lambda}} + e^{\frac{g(z)}{\lambda}} \int_0^z e^{-\frac{g(\zeta)}{\lambda}} h'(\zeta) d\zeta. \quad (1)$$

Therefore, the operator $T_g - \lambda I$ is invertible on a space X precisely when the operator $R_g(\lambda) : X \rightarrow X$ defined above is bounded on X . Of course, this implies the norm equivalence $\|f\|_X \simeq \|R_g(\lambda) f\|_X$ for $f \in X$. This simple but crucial observation has been fruitfully employed in the study of the spectrum in [2], and later in [3] and [11], where connections have been established between the boundedness of $R_g(\lambda)$ and the theory of Muckenhoupt weights and Békollé-Bonami weights. These connections will be our principal tools when establishing the main results of this paper.

In the Hardy space setting the boundedness of the operator $R_g(\lambda)$ is equivalent to a certain weight satisfying a Muckenhoupt \mathcal{A}_∞ -condition. A weight for us will be a positive measurable function, and the \mathcal{A}_∞ -class will consist of weights w on \mathbb{T} for which there exists a constant $C > 0$ such that the following estimate holds for all arcs $I \subseteq \mathbb{T}$:

$$\frac{1}{m(I)} \int_I w dm \leq C \exp \left(\frac{1}{m(I)} \int_I \log w dm \right). \quad (2)$$

Here $m(I)$ denotes the normalized Lebesgue measure of the arc I . We will also need to consider the Muckenhoupt \mathcal{A}_2 -class, i.e the weights satisfying

$$\sup_{I \subseteq \mathbb{T}} \left(\frac{1}{m(I)} \int_I w \, dm \right) \left(\frac{1}{m(I)} \int_I w^{-1} \, dm \right) < \infty. \quad (3)$$

It is straightforward to check that $w \in \mathcal{A}_2$ if and only if $w, w^{-1} \in \mathcal{A}_\infty$. By the Hölder inequality, we can easily prove that \mathcal{A}_∞ is closed under log-convex combinations, in the sense that whenever $w_1, w_2 \in \mathcal{A}_\infty$ and $0 < r < 1$, then $w_1^r w_2^{1-r} \in \mathcal{A}_\infty$. See [10, Chapter 9] for more details.

The following result by Aleman and Peláez characterizes the resolvent set $\rho(T_g|H^p)$ in terms of (2).

Theorem 1.1 (Theorem C of [3]). *Assume that $\lambda \in \mathbb{C} \setminus \{0\}$, $0 < p < \infty$ and $g \in BMOA$. Then, the following assertions are equivalent:*

- (i) $\lambda \in \rho(T_g|H^p)$,
- (ii) the weight $\exp(p \operatorname{Re}(g(e^{it})/\lambda))$ satisfies the \mathcal{A}_∞ -condition.

A similar characterization has been obtained by Aleman and Constantin in the Bergman space setting, and was further refined by Aleman, Pott and Reguera in [5]. There the Békollé-Bonami weights appear instead, which are the Bergman space counterparts of the Muckenhoupt weights. The Békollé-Bonami class B_2 consists of weights on \mathbb{D} satisfying an analogue of (3):

$$\sup_{I \subseteq \mathbb{T}} \left(\frac{1}{A(S_I)} \int_{S_I} w \, dA \right) \left(\frac{1}{A(S_I)} \int_{S_I} w^{-1} \, dA \right) < \infty, \quad (4)$$

where S_I denotes the usual Carleson square associated to the arc $I \subseteq \mathbb{T}$:

$$S_I = \left\{ z \in \mathbb{D} : \frac{z}{|z|} \in I, 1 - m(I) < |z| < 1 \right\},$$

and $A(S_I)$ is the normalized area measure of S_I . There is an analogue of the \mathcal{A}_∞ -class for the Békollé-Bonami weights which is often denoted by B_∞ . The class consists of weights on \mathbb{D} which satisfy

$$\sup_{I \subseteq \mathbb{T}} \left(\int_{S_I} w \, dA \right)^{-1} \left(\int_{S_I} M(w|_{S_I}) \, dA \right) < \infty,$$

where M denotes the usual Hardy-Littlewood maximal function over Carleson squares, and $w|_{S_I}$ is the restriction of w to the square S_I . The weights appearing in our context

will be of the form $w(z) = |e^{g(z)}|$ with $g \in \mathcal{B}$, and in the case of such weights the results of [5] show that the B_∞ -class can be characterized in a similar way to (2), i.e for all arcs $I \subseteq \mathbb{T}$:

$$\frac{1}{A(S_I)} \int_{S_I} w \, dA \leq C \exp \left(\frac{1}{A(S_I)} \int_{S_I} \log w \, dA \right). \quad (5)$$

Again, similarly to the Muckenhoupt weights we have that the weights w, w^{-1} satisfy (5) if and only if $w \in B_2$. Moreover, the weights satisfying (5) are closed under log-convex combinations.

The condition that appears in the characterization of the spectrum of T_g operators acting on the weighted Bergman spaces is the following.

Theorem 1.2 (Part of Theorem A of [2], Corollary 4.6 of [5]). *Let $p > 0$, $\alpha > -1$ and $g \in \mathcal{B}$. For $\lambda \in \mathbb{C} \setminus \{0\}$ the following are equivalent:*

- (i) $\lambda \in \rho(T_g|L_a^{p,\alpha})$,
- (ii) *the weight $(1 - |z|^2)^\alpha \exp(p \operatorname{Re}(g(z)/\lambda))$ satisfies the B_∞ -condition.*

We will use the form of the resolvent in (i) and the conditions of Theorem 1.1 and Theorem 1.2 in the proofs of our main results. The results are stated in Section 2, together with a discussion. The proofs are deferred to Section 3.

2 MAIN RESULTS

Our first main result is a spectral stability property. This is a version of a result which appears in context of so-called growth classes in [II, Theorem 5.4].

Theorem 2.1. *Let g, h be analytic functions such that $T_g, T_h : X \rightarrow X$ are bounded, where $X = H^p$ or $X = L_a^{p,\alpha}$ and $0 < p < \infty$. Suppose that the spectrum $\sigma(T_h|X)$ equals $\{0\}$. Then*

$$\sigma(T_{g+h}|X) = \sigma(T_g|X).$$

The above stability property has previously been studied and established only in very special cases, for instance when g' is a rational function, and h is bounded or induces a compact operator T_h . See, for instance, [13, Theorem 5.2], [2, Theorem B] or [1, Theorem 2.6]. All these cases are covered by Theorem 2.1.

Our next result extends the applicability of Theorem 2.1 by identifying a large class of quasi-nilpotent T_g operators. The result holds true in a much larger class of spaces than just the Hardy and Bergman spaces.

Theorem 2.2. *Let X be a Banach space of analytic functions in \mathbb{D} which contains the constants and such that the algebra $B(X)$ of bounded linear operators on X contains the multiplication operators M_h and the generalized Cesàro operators T_h whenever $h \in H^\infty$. Then we have that $\sigma(T_g|X) = \{0\}$ whenever T_g lies in the norm-closure of $\{T_h : h \in H^\infty\}$ in $B(X)$.*

The conclusion of Theorem 2.2 also holds in the case of the metric spaces H^p and $L_a^{p,\alpha}$ for $p \in (0, 1)$, as will be clear from the proof given in Section 3. The result is particularly useful in case that the space of symbols inducing bounded operators is known, as is the case for the Hardy and Bergman spaces. For instance, the following consequence is immediate from the well-known norm comparabilities $\|T_g\|_{H^p} \simeq \|g\|_{\text{BMOA}}$ and $\|T_g\|_{L_a^{p,\alpha}} \simeq \|g\|_{\mathcal{B}}$.

Corollary 2.3. *If g lies in the norm-closure of H^∞ in BMOA or in the norm-closure of H^∞ in \mathcal{B} , then we have that $\sigma(T_g|H^p) = \{0\}$ and that $\sigma(T_g|L_a^{p,\alpha}) = \{0\}$, respectively.*

It is natural to ask what can be said about the converse statement. If $g \in \text{BMOA}$ or $g \in \mathcal{B}$, and the operator T_g is quasi-nilpotent on H^p or on $L_a^{p,\alpha}$, is then g necessarily contained in the closure of H^∞ in BMOA , or \mathcal{B} , respectively? The Bergman case is related to a long-standing open problem which will be discussed below. In the case of the Hardy spaces the converse does hold. In fact we prove a stronger statement, with a nice geometric flavour.

Theorem 2.4. *Let $0 < p < \infty$ and $g \in \text{BMOA}$. If the spectrum $\sigma(T_g|H^p)$ does not contain any non-zero points of the real and imaginary axes, then g lies in the norm-closure of H^∞ in BMOA , and thus $\sigma(T_g|H^p) = \{0\}$.*

In the statement of Theorem 2.4 the real and imaginary axes can be replaced by two arbitrary lines which intersect orthogonally at the origin. To the authors' best knowledge, all the explicit computations of spectra of T_g on the Hardy spaces, for particular symbols g , reveal that if the spectrum is bigger than $\{0\}$, then it contains a disk with the origin on its boundary (see, for instance, [4]). This obviously implies that there are non-zero points in the spectrum which lie on either the real or imaginary axis.

Theorem 1.1 and Theorem 1.2 provide a deep link between the structure of the spectrum and the theory of exponential weights. In our last main result, we show that the log-convex property of the weight classes implies further properties of the spectrum.

Theorem 2.5. *Let $X = H^p$ or $X = L_a^{p,\alpha}$ for some $0 < p < \infty$, and $T_g : X \rightarrow X$ be bounded. For every non-zero $\lambda \in \sigma(T_g|X)$ and every $0 < r < 1$ there exists a circular arc $I_{r,\lambda}$ centered at $r\lambda$ (see Figure 1), such that the circular sector $S_{r,\lambda}$, created by taking the convex hull of the origin and $I_{r,\lambda}$ is contained in $\sigma(T_g|X)$.*

The proof of Theorem 2.4 ultimately relies on distance formulas of Jones and Garnett which are a consequence of the well-known equivalence of the \mathcal{A}_2 -condition and the Helson-Szegő condition (see Chapter VI of [9]). A similar equivalence is not known in the case of Békollé-Bonami class B_2 . In particular the missing link is the rightmost inequality

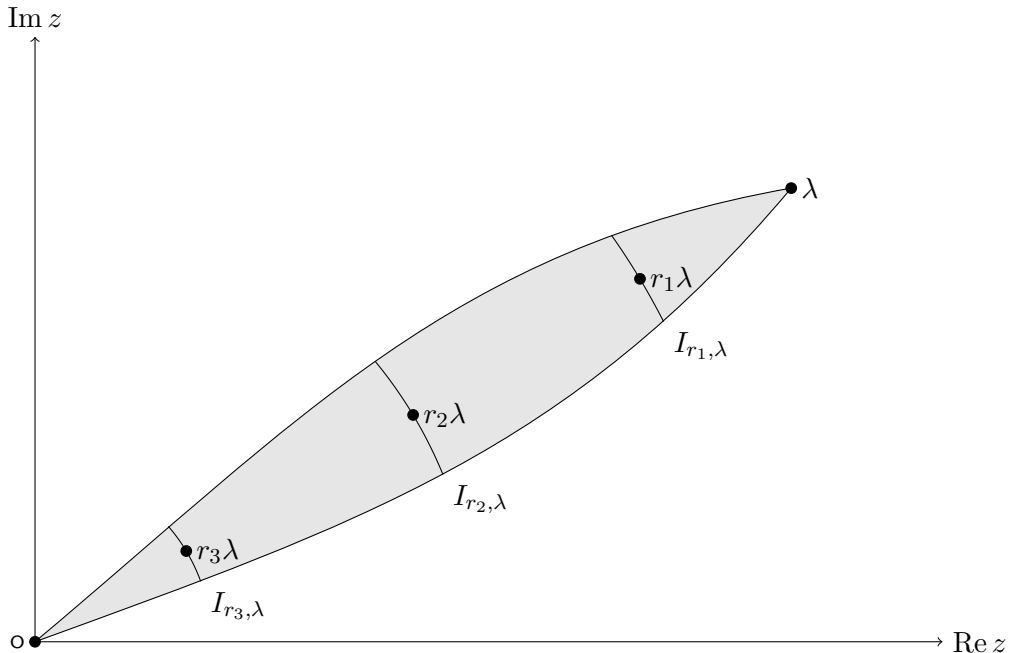


Figure 1: Points in the shaded area belong to the spectrum.

of (10) and consequently the problem of establishing the converse of Corollary 2.3 in the case of the Bergman spaces remains unsolved.

Conjecture. Let $g \in \mathcal{B}$ be such that the weights $\exp(\operatorname{Re}(g(z)/\lambda))$ satisfy (5) for all $\lambda \in \mathbb{C} \setminus \{0\}$. Then g lies in the closure of H^∞ in the Bloch norm.

With the characterization of Theorem 1.2 in mind, the conjecture asserts that the quasinilpotency of T_g on the Bergman spaces is equivalent to g lying in the closure of H^∞ in the Bloch space. The problem of characterizing this closure has been first stated in [12] and remains open to this date. However, there exist results by Galán and Nicolau [7] where a characterization of the closure of the H^p -spaces in the Bloch space is obtained in terms of square-type functions. This characterization does unfortunately not extend to H^∞ , as was shown by a counterexample in [8].

3 PROOFS

Proof of Theorem 2.1. We treat the case when T_g and T_h act on the H^p -spaces. We shall prove the inclusion $\rho(T_g|H^p) \subseteq \rho(T_{g+h}|H^p)$. The reverse inclusion follows by considering $g + h$ and $-h$ instead of g and h .

Fix $\lambda \in \rho(T_g|H^p)$. Then by Theorem 1.1 the weight $\exp(p \operatorname{Re}(g(e^{it})/\lambda))$ satisfies the condition (2), and since the set $\rho(T_g|H^p)$ is open, the same holds for the weight

$$w_1(e^{it}) = \exp(pp' \operatorname{Re}(g(e^{it})/\lambda))$$

whenever $p' > 1$ is sufficiently close to 1. On the other hand, the weight

$$w_2(e^{it}) = \exp(pq' \operatorname{Re}(h(e^{it})/\lambda))$$

satisfies the condition (2) for all real q' by Theorem 1.1 and the assumption that $\sigma(T_h|H^p) = \{0\}$. Let us now fix $p', q' > 1$ such that $1/p' + 1/q' = 1$ and so that w_1 satisfies (2). Let

$$w(e^{it}) = \exp(p \operatorname{Re}((g(e^{it}) + h(e^{it}))/\lambda)).$$

Note that $w = w_1^{1/p'} w_2^{1/q'}$. To show that $\lambda \in \rho(T_{g+h}|H^p)$ we now verify (2) for w :

$$\begin{aligned} \frac{1}{m(I)} \int_I w \, dm &\leq \left(\frac{1}{m(I)} \int_I w_1 \, dm \right)^{1/p'} \left(\frac{1}{m(I)} \int_I w_2 \, dm \right)^{1/q'} \\ &\leq C \exp \left(\frac{1}{m(I)} \int_I (1/p') \log w_1 \, dm \right) \cdot \exp \left(\frac{1}{m(I)} \int_I (1/q') \log w_2 \, dm \right) \\ &= C \exp \left(\frac{1}{m(I)} \int_I (1/p') \log w_1 + (1/q') \log w_2 \, dm \right) \\ &= C \exp \left(\frac{1}{m(I)} \int_I \log w \, dm \right). \end{aligned}$$

This proves the theorem for the case $X = H^p$. The case $X = L_a^{p,\alpha}$ is treated in an analogous way, by using instead the characterization of Theorem 1.2, the openness of the resolvent set and Hölder's inequality. We leave out the details of the computation which are similar to the above one. \square

Proof of Theorem 2.2. Fix $\lambda \in \mathbb{C} \setminus \{0\}$. Pick $h \in H^\infty$ with $\|T_g - T_h\|_{B(X)} = \|T_{g-h}\|_{B(X)}$ small enough for the operator $T_{g-h} - \lambda I$ to be invertible. For the constant function $1 \in X$ we have

$$R_g(\lambda)1 = e^{\frac{h}{\lambda}} e^{\frac{g-h}{\lambda}} = M_{e^{h/\lambda}} R_{g-h}(\lambda)1 \in X.$$

On the other hand, if $f \in X$ with $f(0) = 0$, then

$$\begin{aligned} R_g(\lambda)f(z) &= e^{\frac{g(z)}{\lambda}} \int_0^z e^{-\frac{g(\zeta)}{\lambda}} f'(\zeta) \, d\zeta \\ &= e^{\frac{h(z)}{\lambda}} e^{\frac{g(z)-h(z)}{\lambda}} \int_0^z e^{-\frac{g(\zeta)-h(\zeta)}{\lambda}} e^{-\frac{h(\zeta)}{\lambda}} f'(\zeta) \, d\zeta \\ &= e^{\frac{h(z)}{\lambda}} e^{\frac{g(z)-h(z)}{\lambda}} \int_0^z e^{-\frac{g(\zeta)-h(\zeta)}{\lambda}} \left[(e^{-\frac{h(\zeta)}{\lambda}} f(\zeta))' - e^{-\frac{h(\zeta)}{\lambda}} \frac{h'(\zeta)}{\lambda} f(\zeta) \right] \, d\zeta \\ &= M_{e^{h/\lambda}} R_{g-h}(\lambda) M_{e^{-h/\lambda}} f(z) + M_{e^{h/\lambda}} R_{g-h}(\lambda) T_{e^{-h/\lambda}} f(z). \end{aligned}$$

Note that since $h \in H^\infty$, also $e^{\pm h/\lambda} \in H^\infty$ and thus the last line above is a sum of compositions of bounded operators on X . The operator $R_g(\lambda)$ is therefore itself bounded on X . \square

In order to prove Theorem 2.4, we will need a lemma which establishes the comparability of the distance to H^∞ of a **BMOA**-function in **BMO**-norm to the distance to $L^\infty(\mathbb{T})$ in **BMO**-norm. According to [9], such a result seems to originate back to the work of D. Sarason in [14]. However, we have not been able to find an explicit proof of this result. For the sake of being self-contained, we include a short proof.

Lemma 3.1. *There exists $C > 0$, such that whenever $g \in \mathbf{BMOA}$,*

$$\inf_{f \in H^\infty} \|g - f\|_{\mathbf{BMO}} \leq C \inf_{h \in L^\infty} \|g - h\|_{\mathbf{BMO}} \quad (6)$$

Proof. Without loss of generality, we may assume that $g(0) = 0$. Now let $\{h_n\}_{n \in \mathbb{N}} \subset L^\infty(\mathbb{T})$ be a sequence which minimizes the **BMO**-distance from g to $L^\infty(\mathbb{T})$. By simple triangle inequality, we have for all $n \in \mathbb{N}$

$$\inf_{f \in H^\infty} \|g - f\|_{\mathbf{BMO}} \leq \|g - h_n\|_{\mathbf{BMO}} + \inf_{f \in H^\infty} \|h_n - f\|_{\mathbf{BMO}}. \quad (7)$$

To estimate the second norm, we use the straightforward continuous embedding $L^\infty(\mathbb{T}) \hookrightarrow \mathbf{BMO}$ and the relation $(H_0^1)^\perp = H^\infty$ to get

$$\inf_{f \in H^\infty} \|h_n - f\|_{\mathbf{BMO}} \leq 2 \inf_{f \in H^\infty} \|h_n - f\|_\infty = 2 \sup_{\substack{F \in H_0^1 \\ \|F\|_{H^1} \leq 1}} \left| \int_{\mathbb{T}} F h_n dm \right|. \quad (8)$$

Since the Szegő projection P is self-adjoint and $P(\bar{g}) = 0$, we can write

$$\int_{\mathbb{T}} F h_n dm = \lim_{r \rightarrow 1^-} \int_{\mathbb{T}} P(F)(r\zeta) h_n(\zeta) dm(\zeta) = \lim_{r \rightarrow 1^-} \int_{\mathbb{T}} F(\zeta) \overline{P(\bar{h}_n - \bar{g})(r\zeta)} dm(\zeta). \quad (9)$$

Now using duality on (9) and the fact that $P : \mathbf{BMO} \rightarrow \mathbf{BMOA}$, we obtain

$$\sup_{\substack{F \in H_0^1 \\ \|F\|_{H^1} \leq 1}} \left| \int_{\mathbb{T}} F h_n dm \right| \leq c \left\| P(\bar{h}_n - \bar{g}) \right\|_{\mathbf{BMO}} \leq c' \|h_n - g\|_{\mathbf{BMO}}.$$

Hence according to (8), we have established that $\inf_{f \in H^\infty} \|h_n - f\|_{\mathbf{BMO}} \lesssim \|h_n - g\|_{\mathbf{BMO}}$, thus going back to (7), we ultimately arrive at

$$\inf_{f \in H^\infty} \|g - f\|_{\mathbf{BMO}} \leq C \|g - h_n\|_{\mathbf{BMO}}.$$

Letting $n \rightarrow \infty$ finishes the proof of the lemma. \square

Proof of Theorem 2.4. Fix $0 < p < \infty$ and suppose we have an analytic function $g \in \mathbf{BMO}$, with the property that $\sigma(T_g|H^p) \cap (\mathbb{R} \cup i\mathbb{R}) = \{0\}$. Set $w_1 := \exp(\operatorname{Re}(g))$ and $w_2 := \exp(\operatorname{Im}(g))$. Now since $\lambda, i\lambda \in \rho(T_g|H^p)$ for all $\lambda \in \mathbb{R} \setminus \{0\}$, Theorem 1.1 yields that the weights $w_1^{1/\lambda}, w_2^{1/\lambda}$ both satisfy the \mathcal{A}_∞ -condition for all $\lambda \in \mathbb{R} \setminus \{0\}$. In particular, since both the weights $w_j^{1/\lambda}, w_j^{-1/\lambda}$ satisfy the \mathcal{A}_∞ -condition for $j = 1, 2$, the weights $w_1^{1/\lambda}, w_2^{1/\lambda}$ are in fact \mathcal{A}_2 -weights, for all $\lambda > 0$. A well-known consequence of the Helson-Szegö theorem (see [9, Chapter VI, Section 6]) is that there exists a constant $c > 0$ such that for all real-valued $\phi \in \mathbf{BMO}$ we have

$$\frac{1}{c} \lambda(\phi) \leq \inf_{h \in L^\infty} \|\phi - h\|_{\mathbf{BMO}} \leq c\lambda(\phi) \quad (10)$$

where $\lambda(\phi) = \inf \left\{ \lambda > 0 : e^{\phi/\lambda} \in \mathcal{A}_2 \right\}$. Applying this fact to $\operatorname{Re}(g), \operatorname{Im}(g) \in \mathbf{BMO}$, we get that $\inf_{h \in L^\infty} \|\operatorname{Re}(g) - h\|_{\mathbf{BMO}} = \inf_{h \in L^\infty} \|\operatorname{Im}(g) - h\|_{\mathbf{BMO}} = 0$, thus $\inf_{h \in L^\infty} \|g - h\|_{\mathbf{BMO}} = 0$. According to Lemma 3.1, this is enough to conclude the proof. \square

Before moving on to the proof of Theorem 2.5 we establish a preliminary lemma.

Lemma 3.2. *Let $X = H^p$ or $X = L_a^{p,\alpha}$ for some $0 < p < \infty$. Then the spectrum $\sigma(T_g|X)$ is star-shaped with respect to the origin, that is, if $\lambda \in \sigma(T_g|X)$, then $r\lambda \in \sigma(T_g|X)$ for each $r \in [0, 1]$. In particular, the spectrum is always simply connected.*

Proof. Without loss of generality we may assume that $\sigma(T_g|X) \setminus \{0\}$ is non-trivial. Now suppose that there exists $0 < r < 1$ and a non-zero point $\lambda_0 \in \sigma(T_g|X)$, such that $r\lambda_0 \in \rho(T_g|X)$. We will show that this implies that λ_0 is not in the spectrum, and thus the claim of the lemma will be established. We shall carry out the proof in the case $X = L_a^{p,\alpha}$. According to Theorem 1.2 it suffices to establish that the weight $w_\alpha(z) = v_\alpha(z) \exp(p \operatorname{Re}(g(z)/\lambda_0))$ satisfies the B_∞ -condition, where $v_\alpha(z) = (1 - |z|^2)^\alpha$. We assumed that $v_\alpha \exp(p \operatorname{Re}(g/r\lambda_0))$ satisfies the B_∞ -condition, hence applying Hölder's inequality followed by the B_∞ -condition, we obtain

$$\begin{aligned} & \frac{1}{A(S_I)} \int_{S_I} w_\alpha dA \\ & \leq \left(\frac{1}{A(S_I)} \int_{S_I} v_\alpha \exp \left(p \operatorname{Re} \left(\frac{g}{r\lambda_0} \right) \right) dA \right)^r \left(\frac{1}{A(S_I)} \int_{S_I} v_\alpha dA \right)^{1-r} \\ & \lesssim \exp \left(\frac{1}{A(S_I)} \int_{S_I} \log \left[v_\alpha \exp \left(p \operatorname{Re} \left(\frac{g}{r\lambda_0} \right) \right) \right] dA \right)^r \left(\frac{1}{A(S_I)} \int_{S_I} v_\alpha dA \right)^{1-r} \end{aligned}$$

$$\begin{aligned} &\lesssim \exp \left(\frac{1}{A(S_I)} \int_{S_I} \log \left[v_\alpha \exp \left(p \operatorname{Re} \left(\frac{g}{r\lambda_0} \right) \right) \right] dA \right)^r \\ &\times \exp \left(\frac{1}{A(S_I)} \int_{S_I} \log(v_\alpha) dA \right)^{1-r} = \exp \left(\frac{1}{A(S_I)} \int_{S_I} \log(w_\alpha) dA \right). \end{aligned}$$

In the second last step we also used the B_∞ -condition on the standard weights v_α , which follows from the simple fact that the Bergman projection P_α is bounded on $L^{p,\alpha}$. Hence by Theorem 1.2, we conclude that $\lambda_0 \in \rho(T_g|L_a^{p,\alpha})$, which contradicts our initial assumption, thus the theorem is proved in the case $X = L_a^{p,\alpha}$. The case $X = H^p$ is more straightforward and treated analogously, using instead the characterization of Theorem 1.1. \square

Proof of Theorem 2.5. This time, we carry out the proof in the case $X = H^p$, where the Bergman case is similar. Pick $\lambda \in \sigma(T_g|H^p) \setminus \{0\}$. By means of multiplying g with a unimodular constant, which corresponds to rotating the spectrum, we may without loss of generality assume that $\lambda > 0$. Now fix $0 < r < r' < 1$ and consider the circular arc parametrization $\gamma_r(t) = r\lambda e^{it}$, $-\pi < t \leq \pi$. Notice that

$$\operatorname{Re} \left(\frac{g}{\gamma_r(t)} \right) = \frac{\operatorname{Re}(g)}{r\lambda} \cos(t) + \frac{\operatorname{Im}(g)}{r\lambda} \sin(t) = \frac{\operatorname{Re}(g)}{\frac{r}{\cos(t)} \cdot \lambda} + \frac{\operatorname{Im}(g)}{\frac{r}{\sin(t)} \cdot \lambda} \quad (\text{II})$$

Let $|\sigma(T_g|H^p)|$ denote the spectral radius of the operator T_g on H^p . Denote by $J_{r,\lambda}$ the interval consisting of t for which $\frac{r}{\cos(t)} \leq r'$ and $\frac{r(1-r')\lambda}{r'|\sin(t)|} > |\sigma(T_g|H^p)|$, i.e.

$$|t| \leq \min \left\{ \arccos(r/r'), \arcsin(r(1-r')\lambda/r'|\sigma(T_g|H^p)|) \right\}.$$

Then by Theorem 1.1 and the star-shaped property of the spectrum, we have that the weight

$$u_t^{r'} := \exp \left(\frac{\operatorname{Re}(g)}{\frac{r}{r' \cos(t)} \cdot \lambda} \right) \notin \mathcal{A}_\infty \quad \forall t \in J_{r,\lambda}, \quad (\text{I2})$$

while $v_t^{r'/(1-r')} := \exp \left(\frac{\operatorname{Im}(g)}{\frac{r(1-r')\lambda}{r' \sin(t)}} \right) \in \mathcal{A}_\infty$, for all $t \in J_{r,\lambda}$. Since the interval $J_{r,\lambda}$ is symmetric around zero and the sine-function is odd, we have that $v_{-t} = v_t^{-1}$, thus $v_t^{-r'/(1-r')}$ is also an \mathcal{A}_∞ -weight. Setting $w_t = \exp \left(\operatorname{Re} \left(\frac{g}{\gamma_r(t)} \right) \right)$ we can rewrite $v_t^{-1} w_t = u_t$. For the sake of obtaining a contradiction, suppose $w_{t_0} \in \mathcal{A}_\infty$, for some $t_0 \in J_{r,\lambda}$, which by Theorem 1.1 corresponds to the assumption that $\gamma_r(t_0) \in \rho(T_g|H^p)$. Now since

$v_{t_0}^{-r'/(1-r')}$ is an \mathcal{A}_∞ -weight and the \mathcal{A}_∞ -class is closed under log-convex combinations, we get that

$$u_{t_0}^{r'} = v_{t_0}^{-r'} w_{t_0}^{r'} = \left(v_{t_0}^{-r'/(1-r')} \right)^{1-r'} w_{t_0}^{r'} \in \mathcal{A}_\infty. \quad (13)$$

This contradicts our initial assumption in (12). We have thus established that the circular arc corresponding to $J_{r,\lambda}$, defined by $I_{r,\lambda} := \{r\lambda e^{it} : t \in J_{r,\lambda}\}$, is contained in $\sigma(T_g|H^p)$. Taking convex hulls of the origin and $I_{r,\lambda}$, we conclude by Lemma 3.2 that the corresponding circular sector $S_{r,\lambda}$ is included in the spectrum. \square

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